

# EUR 2584.e

EUROPEAN ATOMIC ENERGY COMMUNITY - EURATOM

A PROGRAMME FOR THE SOLUTION OF THE  
MONOENERGETICAL TRANSPORT EQUATION  
IN THE P5 APPROXIMATION  
FOR A MULTIREGION CYLINDRICAL GEOMETRY

by

L. MASSIMO

1965



Joint Nuclear Research Center  
Ispra Establishment - Italy

Reactor Physics Department  
Reactor Theory and Analysis



## LEGAL NOTICE

This document was prepared under the sponsorship of the Commission of the European Atomic Energy Community (EURATOM).

Neither the EURATOM Commission, its contractors nor any person acting on their behalf:

Make any warranty or representation, express or implied, with respect to the accuracy, completeness, or usefulness of the information contained in this document, or that the use of any information, apparatus, method, or process disclosed in this document may not infringe privately owned rights; or

Assume any liability with respect to the use of, or for damages resulting from the use of any information, apparatus, method or process disclosed in this document.

This report is on sale at the addresses listed on cover page 4

at the price of FF 5,—	FB 50,—	DM 4,—	Lit. 620	Fl. 3,60
------------------------	---------	--------	----------	----------

**When ordering, please quote the EUR number and the title, which are indicated on the cover of each report.**

Printed by Guyot, s.a.  
Brussels, November 1965

This document was reproduced on the basis of the best available copy.



## EUR 2584.e

A PROGRAMME FOR THE SOLUTION OF THE MONOENERGETICAL TRANSPORT EQUATION IN THE P5 APPROXIMATION FOR A MULTIGROUP CYLINDRICAL GEOMETRY by L. MASSIMO

European Atomic Energy Community - EURATOM  
Joint Nuclear Research Center - Ispra Establishment (Italy)  
Reactor Physics Department  
Reactor Theory and Analysis  
Brussels, November 1965 - 40 pages - FB 50

The report describes a Fortran II code for the IBM 7090 computer which solves the monoenergetical transport equation in the P5 approximation for a multiregion cylindrical geometry.

Special attention is paid to the numerical difficulties encountered in this work.

---

## EUR 2584.e

A PROGRAMME FOR THE SOLUTION OF THE MONOENERGETICAL TRANSPORT EQUATION IN THE P5 APPROXIMATION FOR A MULTIGROUP CYLINDRICAL GEOMETRY by L. MASSIMO

European Atomic Energy Community - EURATOM  
Joint Nuclear Research Center - Ispra Establishment (Italy)  
Reactor Physics Department  
Reactor Theory and Analysis  
Brussels, November 1965 - 40 pages - FB 50

The report describes a Fortran II code for the IBM 7090 computer which solves the monoenergetical transport equation in the P5 approximation for a multiregion cylindrical geometry.

Special attention is paid to the numerical difficulties encountered in this work.

---

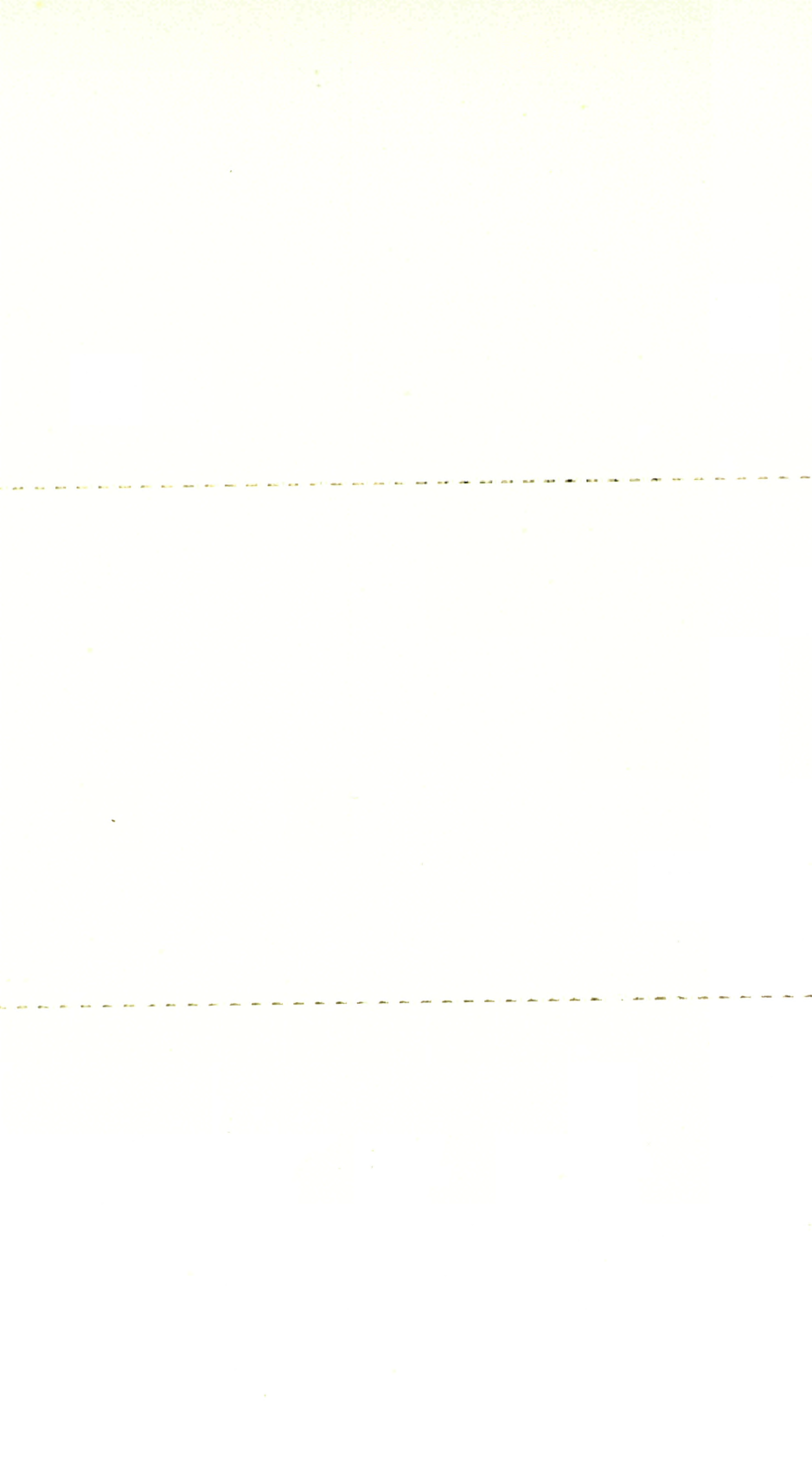
## EUR 2584.e

A PROGRAMME FOR THE SOLUTION OF THE MONOENERGETICAL TRANSPORT EQUATION IN THE P5 APPROXIMATION FOR A MULTIGROUP CYLINDRICAL GEOMETRY by L. MASSIMO

European Atomic Energy Community - EURATOM  
Joint Nuclear Research Center - Ispra Establishment (Italy)  
Reactor Physics Department  
Reactor Theory and Analysis  
Brussels, November 1965 - 40 pages - FB 50

The report describes a Fortran II code for the IBM 7090 computer which solves the monoenergetical transport equation in the P5 approximation for a multiregion cylindrical geometry.

Special attention is paid to the numerical difficulties encountered in this work.



**EUR 2584.e**

EUROPEAN ATOMIC ENERGY COMMUNITY - EURATOM

A PROGRAMME FOR THE SOLUTION OF THE  
MONOENERGETICAL TRANSPORT EQUATION  
IN THE P5 APPROXIMATION  
FOR A MULTIREGION CYLINDRICAL GEOMETRY

by

L. MASSIMO

1965



Joint Nuclear Research Center  
Ispra Establishment - Italy

Reactor Physics Department  
Reactor Theory and Analysis

describes a Fortran II code for the IBM 7090 computer which  
computes the exact solution in the TD approximation for

to this

## Contents

	Page
1. General description	"
2. Theory and Code	"
3. References	"
4. Acknowledgment	"
Appendix 1 - The air-gap conditions	"
Appendix 2 - Numerical calculations of the $G(\gamma)$ functions	"
Appendix 3 - Input description	"
Appendix 4 - Code list and sample problem	"

## SPHERICAL HARMONICS P5 PROGRAMME<sup>+</sup>

### 1. General description

The code described in this report was written by the author while at the Dragon Project, A.E.E. Winfrith, for the Ferranti Mercury computer.

Because of the difficulty in the use of double precision matrix operations on that computer the code did not work properly there for many difficult cases. Now the programme has been translated into Fortran II for the IBM 7090 computer and it works satisfactorily.

The programme solves the monoenergetical transport equation in the  $P_5$  approximation for a multiregion cylindrical geometry. Up to 10 concentric regions can be considered. The scattering is considered to be isotropic.

Two boundary conditions, reflective and infinite medium, are available, and an air gap may be considered. The equations are solved by the propagation method described for the  $P_3$  approximation in ref. 1.

Such methods have the advantage of solving large regions without losing accuracy in cases where numerical techniques might fail, and are quick where large uniform regions are considered.

While the programming of the  $P_3$  did not present great numerical difficulties, the  $P_5$  required double precision and a tightening up of the procedure of matrix propagation.

Being all analytical the code is very fast (a typical case can run in 15 to 20 seconds) but the numerical difficulties encountered in writing it show that the  $P_5$  is about the highest approximation which one can treat with this method.

As a control on the accuracy of the results, the flux at

---

<sup>+</sup> Manuscript received on October 5, 1965



each boundary is printed twice with 8 decimal figures, once calculated from the inner side and the other time from the outer side of the boundary. Any difference between the two fluxes is due to numerical inaccuracy.

## 2. Theory and Code

The expansion of the flux is given by

$$\begin{aligned}
 2\pi\psi(r, \varphi) = & \\
 & \frac{1}{2} \psi_{00}(r) + \frac{3}{2} \psi_{11}(r) P_1^1(\eta) \cos\omega + 5 \left[ \frac{1}{2} P_2(\eta) \psi_{02}(r) + \right. \\
 & \left. + \frac{1}{4!} P_2^2(\eta) \cos 2\omega \psi_{22}(r) \right] + 7 \left[ \frac{2}{4!} P_3^1(\eta) \cos\omega \psi_{13}(r) + \right. \\
 & \left. + \frac{1}{6!} P_3^3(\eta) \cos 3\omega \psi_{33}(r) \right] + \\
 & + 9 \left[ \frac{1}{2} P_4(\eta) \psi_{04}(r) + \frac{2}{6!} P_4^2(\eta) \cos 2\omega \psi_{24}(r) + \right. \\
 & \left. + \frac{1}{8!} P_4^4(\eta) \cos 4\omega \psi_{44}(r) \right] + 11 \left[ \frac{4}{6!} P_5^1(\eta) \cos\omega \psi_{15}(r) + \right. \\
 & \left. + \frac{2}{8!} P_5^3(\eta) \cos 3\omega \psi_{35}(r) + \frac{1}{10!} P_5^5(\eta) \cos 5\omega \psi_{55}(r) \right] \\
 & \dots\dots\dots (1)
 \end{aligned}$$

and the expressions for the various moments  $\psi_{mn}$  in the P5 approximation are as follows:

$$\psi_{00}(r) = \sum_{i=1}^3 \{B_i I_0(\gamma_i r) + C_i K_0(\gamma_i r)\}$$

$$\psi_{11}(r) = \sum_i G_1(\gamma_i) \{B_i I_1(\gamma_i r) - C_i K_1(\gamma_i r)\}$$

$$\psi_{02}(r) = \sum_i -\frac{1}{2} G_2(\gamma_i) \{B_i I_0(\gamma_i r) + C_i K_0(\gamma_i r)\} +$$

$$+ \sum_{j=1}^2 \{M_j I_0(\beta_j r) + N_j K_0(\beta_j r)\}$$

$$\psi_{22}(r) = \sum_i 3 G_2(\gamma_i) \{B_i I_2(\gamma_i r) + C_i K_2(\gamma_i r)\} +$$

$$+ 2 \sum_j \{M_j I_2(\beta_j r) + N_j K_2(\beta_j r)\}$$

$$\psi_{13}(r) = \sum_i -\frac{3}{2} G_3(\gamma_i) \{B_i I_1(\gamma_i r) - C_i K_1(\gamma_i r)\} -$$

$$- 5 \sum_j \frac{1}{\beta_j} \{M_j I_1(\beta_j r) - N_j K_1(\beta_j r)\}$$

$$\psi_{33}(r) = \sum_i 15 G_3(\gamma_i) \{B_i I_3(\gamma_i r) - C_i K_3(\gamma_i r)\} -$$

$$- 30 \sum_j \frac{1}{\beta_j} \{M_j I_3(\beta_j r) - N_j K_3(\beta_j r)\}$$

$$\psi_{04}(r) = \sum_i \frac{3}{8} G_4(\gamma_i) \{B_i I_0(\gamma_i r) + C_i K_0(\gamma_i r)\} +$$

$$+ \sum_j \frac{5(\beta_j^2 - 7)}{12 \beta_j^2} \{M_j I_0(\beta_j r) + N_j K_0(\beta_j r)\} + Q I_0(ur) + R K_0(ur)$$



$$\begin{aligned}\psi_{24}(r) = & \sum_i -\frac{15}{2} G_4(\gamma_i) \{B_i I_2(\gamma_i r) + C_i K_2(\gamma_i r)\} - \\ & - \sum_j \frac{5(\beta_j^2 - 7)}{\beta_j^2} \{M_j I_2(\beta_j r) + N_j K_2(\beta_j r)\} + 12\{Q I_2(vr) + \\ & + R K_2(vr)\}\end{aligned}$$

$$\begin{aligned}\psi_{44}(r) = & \sum_i 105 G_4(\gamma_i) \{B_i I_4(\gamma_i r) + C_i K_4(\gamma_i r)\} - \\ & - \sum_j \frac{70(\beta_j^2 - 7)}{\beta_j^2} \{M_j I_4(\beta_j r) + N_j K_4(\beta_j r)\} + 24\{Q I_4(vr) + \\ & + R K_4(vr)\}\end{aligned}$$

$$\begin{aligned}\psi_{15}(r) = & \sum_i \frac{15}{8} G_5(\gamma_i) \{B_i I_1(\gamma_i r) - C_i K_1(\gamma_i r)\} - \\ & - \sum_j \frac{35(\beta_j^2 - 7)}{44\beta_j} \{M_j I_1(\beta_j r) - N_j K_1(\beta_j r)\} - \frac{9}{v} \{Q I_1(vr) - \\ & - R K_1(vr)\}\end{aligned}$$

$$\begin{aligned}\psi_{35}(r) = & \sum_i -\frac{105}{2} G_5(\gamma_i) \{B_i I_3(\gamma_i r) - C_i K_3(\gamma_i r)\} + \\ & + \sum_j \frac{105}{11} \frac{(\beta_j^2 - 7)}{\beta_j} \{M_j I_3(\beta_j r) - N_j K_3(\beta_j r)\} - \frac{324}{v} \{Q I_3(vr) - R K_3(vr)\}\end{aligned}$$

$$\begin{aligned} \psi_{55}(r) = & \sum_i 945 G_5(\gamma_i) \{B_i I_5(\gamma_i r) - C_i K_5(\gamma_i r)\} + \\ & + \sum_j \frac{3150}{11} \frac{(\beta_j^2 - 7)}{\beta_j} \{M_j I_5(\beta_j r) - N_j K_5(\beta_j r)\} - \frac{1080}{v} \{Q I_5(vr) - \\ & - R K_5(vr)\} \end{aligned} \quad \dots\dots\dots(2a)$$

where

$$\begin{aligned} G_0 &= 1 & a &= 1 - c \\ G_1 &= \frac{a}{\gamma} & c &= N^\circ \text{ of secondaries per collision} \\ G_{n+1} &= -\frac{1}{n+1} \left\{ \frac{2n+1}{\gamma} G_n + n G_{n-1} \right\} \end{aligned} \quad \dots\dots\dots(2b)$$

and  $\gamma_1, \gamma_2, \gamma_3$  are the positive roots of

$$25\gamma^6 - 21(14 + 11a)\gamma^4 + 35(11 + 34a)\gamma^2 - 1155a = 0$$

(see appendix 2)

$\beta_1$  and  $\beta_2$  are the positive roots of

$$\beta^4 - 18\beta^2 + 33 = 0$$

$$\text{and } v = \sqrt{11} \quad \dots\dots\dots(2c)$$

We can consider the flux as a vector whose components are given by the moments  $\psi_{mn}$  so that at a point  $r$  in a given region  $i$ , it may be represented by

$$\underline{\psi}_i^r = A_i^r \underline{C}_i + \underline{S}_i \quad (3)$$

$A_i^r$  is a  $(12 \times 12)$  matrix of modified Bessel functions whose arguments are function of the region number  $i$  and position  $r$ .

$\underline{S}_i^r$  is the source vector (as the source is supposed to be isotropic all the components are zero but the first one).



$\underline{C}_i$  is the vector of the 12 integration constants [ $B_1, C_1, B_2, \dots, Q, R$ ] to be determined from the boundary conditions imposed.

If  $v$  is the number of regions and the boundaries are  $r_1, r_2, \dots, r_n$ , the interface conditions of continuity will lead to the following set of  $(n - 1)$  equations:

$$A_1^{r_1} \underline{C}_1 + \underline{S}_1 = A_2^{r_1} \underline{C}_2 + \underline{S}_2$$

$$A_2^{r_2} \underline{C}_2 + \underline{S}_2 = A_3^{r_2} \underline{C}_3 + \underline{S}_3$$

$$A_{n-1}^{r_{n-1}} \underline{C}_{n-1} + \underline{S}_{n-1} = A_n^{r_{n-1}} \underline{C}_n + \underline{S}_n$$

by defining  $B_i = (A_{i+1}^{r_i})^{-1}$  we obtain the recurrent formula

$$\underline{C}_{i+1} = B_i (A_i^{r_i} \underline{C}_i + \underline{S}_i - \underline{S}_{i+1}) \quad (4)$$

which propagates the coefficients  $\underline{C}_i$  outwards, region by region.

The coefficients  $\underline{C}_i$  are found by applying the boundary conditions of finiteness of the flux in the centre (the coefficients associated with the K functions must be zero) and the chosen boundary conditions at the last region. If the last region is infinite, in  $\underline{C}_n$  the coefficients associated with the I functions must vanish.

If the last boundary is a reflective boundary the components  $\psi_{11} ; \psi_{13} ; \psi_{33} ; \psi_{15} ; \psi_{35} ; \psi_{55}$  must vanish [2]. In this latter case the condition will not be applied to  $\underline{C}_n$  but to

$$(A_n^{r_n} \underline{C}_n + \underline{S}_n).$$

In both cases we start from the centre with a starter  $\underline{C}_1$  (12 x 1) with 6 zero components and we must determine the other 6 in such a way that the (12 x 1) response  $\underline{R}$  that we get after propagation has zeros in 6 defined positions.

The response  $\underline{R}$  will be the sum of a homogeneous response  $\underline{R}^*$  (that we would get with all  $\underline{S}_i = 0$ ) and an effect  $\underline{b}$  of the sources (that we can get by propagating a  $\underline{C}_1$  of all zeros).

$$\underline{R} = \underline{R}^* + \underline{b}$$

If we propagate a matrix  $C$  of 6 independent starters (each satisfying to the centre boundary condition) as for example:

1	0	0	0	0	0
0	0	0	0	0	0
0	1	0	0	0	0
0	0	0	0	0	0
0	0	1	0	0	0
0	0	0	0	0	0
0	0	0	1	0	0
0	0	0	0	0	0
0	0	0	0	1	0
0	0	0	0	0	0
0	0	0	0	0	1
0	0	0	0	0	0

with all  $\underline{S}_i = 0$ .

We get an homogeneous response  $\underline{R}^*$ .

If  $\underline{\phi}$  is a (6 x 1) vector we will have that

$$\underline{C}_1 \underline{\phi} \rightarrow \underline{R}^* \underline{\phi} + \underline{b}$$

We need to find  $\underline{\phi}$  such that  $\underline{R}^* \underline{\phi} + \underline{b}$  satisfies the outer boundary condition, and  $\underline{C}_1 \underline{\phi}$  will be the starter  $\underline{C}_1$  that solves our problem.

Because of loss of accuracy it is necessary to tighten up the process of obtainment of  $\underline{R}^*$ .

To do so we select from  $\underline{R}^*$  (which is a 12 x 6 matrix) the six rows which are of interest for imposing the boundary conditions (i.e. the rows corresponding to the coefficients of the



I functions in case of infinite medium, or the rows corresponding to the 6 components of the flux which must vanish in the case of reflective boundary conditions).

Let us call  $Q$  this  $6 \times 6$  matrix. If we propagate as a new starter  $C'_1$  the product  $C'_1 Q^{-1}$  we will obtain a new  $R^*$ , out of which we can select a new  $Q_1$ . This new  $Q_1$  should be a unit matrix if there would be no loss of accuracy in the process. If  $Q_1$  is not near enough to a unit matrix the loop is continued repropagating  $C'_1 Q_1^{-1}$  until a good enough unit matrix is obtained.

At the end we obtain a starter which gives a response  $R^*$  which is a unit matrix on the rows of interest for the boundary conditions.

At this point the vector  $\underline{\phi}$  is easily determined because  $R^*\underline{\phi}$  will be a  $(12 \times 1)$  vector having the components of  $\underline{\phi}$  on the positions of interest for the boundary conditions.

Because we want the  $12 \times 1$  vector  $R^*\underline{\phi} + \underline{b}$  to have zeros on those positions, the components of  $\underline{\phi}$  will be equal and of opposite sign of the components of  $\underline{b}$  corresponding to those positions.

For example in the case of infinite medium, calling  $\phi_1, \phi_2, \dots, \phi_6$  the components of  $\underline{\phi}$  we have

$R^* =$

1	0	0	0	0	0
0	1	0	0	0	0
0	0	1	0	0	0
0	0	0	1	0	0
0	0	0	0	1	0
0	0	0	0	0	1

so that

$$R^* \phi =$$

$$\begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \\ \phi_4 \\ \phi_5 \\ \phi_6 \end{pmatrix}$$

and to satisfy the boundary conditions we will need:

$$\begin{aligned} \phi_1 &= -b_1 \\ \phi_2 &= -b_3 \\ \phi_3 &= -b_5 \\ \phi_4 &= -b_7 \\ \phi_5 &= -b_9 \\ \phi_6 &= -b_{11} \end{aligned}$$

At this point the programme can calculate the coefficients  $\underline{C}_i$  in every region. If we call  $C_i^*$  the (12x6) matrices that we got as intermediate results while calculating  $R^*$ , and  $\underline{b}_i$  the corresponding intermediate results obtained while calculating  $\underline{b}$ , we have:

$$\underline{C}_i = C_i^* \underline{\phi} + \underline{b}_i$$

The flux is then calculated in the requested points.

The output also contains the mean flux in every region and an asymptotic flux and current for extrapolation lengths calculations.



### 3. References

1. J.R.Askew, R.J.Brissenden  
Mercury programme 560 - Spherical harmonics  $P_3$  approximation - One group, cylindrical geometry.  
AEEW-M116, January 1961
2. J.Tait  
The calculation of the fine structure of the thermal neutron flux in a pile, by the spherical harmonics method.  
A/CONF. 8/P/433 - 1955

### 4. Acknowledgment

Thanks are due to Mr. R.J.Brissenden of U.K.A.E.A. Winfrith for his help given in overcoming the difficulties encountered while first programming this work.

The translation into Fortran is largely due to Mr. G.Fassone.

# A p p e n d i x 1

## The air-gap condition

The equations connecting the various moments across the air gap are as follows:

$$a \psi_{11}(a) = c \psi_{11}(c) \quad (5.1)$$

$$a \psi_{13}(a) = c \psi_{13}(c) \quad (5.2)$$

$$a \psi_{15}(a) = c \psi_{15}(c) \quad (5.3)$$

$$a^3 \{ \psi_{33}(a) - 12 \psi_{11}(a) + 2 \psi_{13}(a) \} = c^3 \{ \psi_{33}(c) - 12 \psi_{11}(c) + 2 \psi_{13}(c) \} \quad \dots\dots\dots (5.4)$$

$$a^3 \{ \psi_{35}(a) - 12 \psi_{11}(a) - 28 \psi_{13}(a) + 12 \psi_{15}(a) \} = c^3 \{ \psi_{35}(c) - 12 \psi_{11}(c) - 28 \psi_{13}(c) + 12 \psi_{15}(c) \} \quad \dots\dots\dots (5.5)$$

$$a^5 \{ \psi_{55}(a) + 1296 \psi_{11}(a) - 336 \psi_{13}(a) + 48 \psi_{15}(a) - 168 \psi_{33}(a) + 6 \psi_{35}(a) \} = c^5 \{ \psi_{55}(c) + 1296 \psi_{11}(c) - 336 \psi_{13}(c) + 48 \psi_{15}(c) - 168 \psi_{33}(c) + 6 \psi_{35}(c) \} \quad \dots\dots\dots (5.6)$$

$$\begin{aligned} & - \frac{\pi}{4} \psi_{00}(a) + \frac{\pi}{2} \psi_{11}(a) + \frac{5\pi}{32} \psi_{02}(a) - \frac{5\pi}{64} \psi_{22}(a) + \frac{9\pi}{256} \psi_{04}(a) - \\ & - \frac{\pi}{256} \psi_{24}(a) + \frac{\pi}{2048} \psi_{44}(a) = \\ & - \frac{\pi}{4} \psi_{00}(c) + \frac{c}{a} \left\{ \sin^{-1} \frac{a}{c} + \frac{a}{c} \sqrt{1 - \frac{a^2}{c^2}} \right\} \psi_{11}(c) + \frac{5\pi}{32} \psi_{02}(c) - \\ & - \frac{15\pi}{64} \left( 1 - \frac{2}{3} \frac{a^2}{c^2} \right) \psi_{22}(c) + \frac{7}{45} \left( 1 - \frac{a^2}{c^2} \right)^{3/2} \psi_{33}(c) + \frac{9\pi}{256} \psi_{04}(c) - \\ & - \frac{3\pi}{256} \left( 1 - \frac{2}{3} \frac{a^2}{c^2} \right) \psi_{24}(c) - \frac{\pi}{2048} \left( 15 - 40 \frac{a^2}{c^2} + 24 \frac{a^4}{c^4} \right) \psi_{44}(c) + \\ & + \frac{11}{1260} \left( 1 - \frac{a^2}{c^2} \right)^{3/2} \psi_{35}(c) + \frac{11}{4200} \left( 1 - \frac{a^2}{c^2} \right)^{3/2} \left( 1 - \frac{8}{3} \frac{a^2}{c^2} \right) \psi_{55}(c) \\ & \dots\dots\dots (5.7) \end{aligned}$$

$$\begin{aligned}
 & - \frac{3\pi}{32} \psi_{00}(a) - \frac{105\pi}{128} \psi_{02}(a) + \frac{5\pi}{256} \psi_{22}(a) + \frac{\pi}{2} \psi_{13}(a) + \\
 & + \frac{2079\pi}{4096} \psi_{04}(a) - \frac{171\pi}{4096} \psi_{24}(a) - \frac{9\pi}{32768} \psi_{44}(a) = - \frac{3\pi}{32} \psi_{00}(c) - \\
 & - \frac{105\pi}{128} \psi_{02}(c) + \frac{15\pi}{256} \left(1 - \frac{2}{3} \frac{a^2}{c^2}\right) \psi_{22}(c) + \frac{c}{a} \left\{ \sin^{-1} \frac{a}{c} + \right. \\
 & + \left. \frac{a}{c} \sqrt{1 - \frac{a^2}{c^2}} \right\} \psi_{13}(c) - \frac{1}{15} \left(1 - \frac{a^2}{c^2}\right)^{3/2} \psi_{33}(c) + \frac{2079\pi}{4096} \psi_{04}(c) - \\
 & - \frac{513\pi}{4096} \left(1 - \frac{2}{3} \frac{a^2}{c^2}\right) \psi_{24}(c) + \frac{9\pi}{32768} \left(15 - 40 \frac{a^2}{c^2} + 24 \frac{a^4}{c^4}\right) \psi_{44}(c) + \\
 & + \frac{11}{210} \left(1 - \frac{a^2}{c^2}\right)^{3/2} \psi_{35}(c) - \frac{11}{6300} \left(1 - \frac{a^2}{c^2}\right)^{3/2} \left(1 - \frac{8}{3} \frac{a^2}{c^2}\right) \psi_{55}(c) \\
 & \dots\dots(5.8)
 \end{aligned}$$

$$\begin{aligned}
 & - \frac{15\pi}{256} \psi_{00}(a) - \frac{1425\pi}{4096} \psi_{02}(c) + \frac{25\pi}{8192} \psi_{22}(a) - \frac{42255\pi}{32768} \psi_{04}(a) + \\
 & + \frac{495\pi}{32768} \psi_{24}(a) + \frac{9\pi}{262144} \psi_{44}(a) + \frac{\pi}{2} \psi_{15}(a) = - \frac{15\pi}{256} \psi_{00}(c) - \\
 & - \frac{1425\pi}{4096} \psi_{02}(c) + \frac{75\pi}{8192} \left(1 - \frac{2}{3} \frac{a^2}{c^2}\right) \psi_{22}(c) - \frac{42255\pi}{32768} \psi_{04}(c) + \\
 & + \frac{1485\pi}{32768} \left(1 - \frac{2}{3} \frac{a^2}{c^2}\right) \psi_{24}(c) - \frac{7\pi}{262144} \left(15 - 40 \frac{a^2}{c^2} + 24 \frac{a^4}{c^4}\right) \psi_{44}(c) + \\
 & + \frac{c}{a} \left\{ \sin^{-1} \frac{a}{c} + \frac{a}{c} \sqrt{1 - \frac{a^2}{c^2}} \right\} \psi_{15}(c) - \frac{1}{28} \left(1 - \frac{a^2}{c^2}\right)^{3/2} \psi_{35}(c) + \\
 & + \frac{1}{2520} \left(1 - \frac{a^2}{c^2}\right)^{3/2} \left(1 - \frac{8}{3} \frac{a^2}{c^2}\right) \psi_{55}(c) \dots\dots(5.9)
 \end{aligned}$$

$$\begin{aligned}
 & \frac{15\pi}{16} \psi_{00}(a) - \frac{75\pi}{64} \psi_{02}(a) - \frac{225\pi}{128} \psi_{22}(a) + \frac{\pi}{2} \psi_{33}(a) + \frac{405\pi}{2048} \psi_{04}(a) + \\
 & + \frac{135\pi}{2048} \psi_{24}(a) - \frac{675\pi}{16384} \psi_{44}(a) = \frac{15\pi}{16} \psi_{00}(c) + 12 \frac{c}{a} \left(1 - \frac{c^2}{a^2}\right) \\
 & \left\{ \sin^{-1} \frac{a}{c} - \frac{a}{c} \sqrt{1 - \frac{a^2}{c^2}} \right\} \psi_{11}(c) - \frac{75\pi}{64} \psi_{02}(c) +
 \end{aligned}$$



$$\begin{aligned}
 & + \frac{375\pi}{128} \left( \frac{1}{3} - \frac{14}{15} \frac{a^2}{c^2} \right) \psi_{22}(c) - 2 \frac{c}{a} \left( 1 - \frac{c^2}{a^2} \right) \left\{ \sin^{-1} \frac{a}{c} - \frac{a}{c} \sqrt{1 - \frac{a^2}{c^2}} \right\} \psi_{13}(c) + \\
 & + \frac{c^3}{a^3} \left\{ \sin^{-1} \frac{a}{c} - \frac{a}{c} \sqrt{1 - \frac{a^2}{c^2}} \left( 1 + \frac{4}{3} \frac{a^2}{c^2} - \frac{10}{3} \frac{a^4}{c^4} \right) \right\} \psi_{33}(c) + \\
 & + \frac{405\pi}{2048} \psi_{04}(c) - \frac{225\pi}{2048} \left( \frac{1}{3} - \frac{14}{15} \frac{a^2}{c^2} \right) \psi_{24}(c) + \\
 & + \frac{15\pi}{16384} \left( 35 - 392 \frac{a^2}{c^2} + 312 \frac{a^4}{c^4} \right) \psi_{44}(c) - \frac{11}{945} \left( 1 - \frac{a^2}{c^2} \right)^{\frac{1}{2}} \\
 & \left( 1 - 17 \frac{a^2}{c^2} + 16 \frac{a^4}{c^4} \right) \psi_{55}(c) \dots\dots\dots(5.10)
 \end{aligned}$$

$$\begin{aligned}
 & \frac{105\pi}{64} \psi_{00}(a) + \frac{5775\pi}{1024} \psi_{02}(a) - \frac{1575\pi}{2048} \psi_{22}(a) - \frac{82215\pi}{8192} \psi_{04}(a) - \\
 & - \frac{20601\pi}{8192} \psi_{24}(a) + \frac{945\pi}{32768} \psi_{44}(a) + \frac{\pi}{2} \psi_{35}(a) = \frac{105\pi}{64} \psi_{00}(c) + \\
 & + 12 \frac{c}{a} \left( 1 - \frac{c^2}{a^2} \right) \left\{ \sin^{-1} \frac{a}{c} - \frac{a}{c} \sqrt{1 - \frac{a^2}{c^2}} \right\} \psi_{11}(c) + \frac{5775\pi}{1024} \psi_{02}(c) + \\
 & + \frac{175\pi}{2048} \left( 5 - 14 \frac{a^2}{c^2} \right) \psi_{22}(c) + 28 \frac{c}{a} \left( 1 - \frac{c^2}{a^2} \right) \left\{ \sin^{-1} \frac{a}{c} - \right. \\
 & - \left. \frac{a}{c} \sqrt{1 - \frac{a^2}{c^2}} \right\} \psi_{13}(c) - \frac{82215\pi}{8192} \psi_{04}(c) + \frac{2289\pi}{8192} \left( 5 - 14 \frac{a^2}{c^2} \right) \psi_{24}(c) - \\
 & - \frac{21\pi}{65536} \left( 35 - 390 \frac{a^2}{c^2} + 312 \frac{a^4}{c^4} \right) \psi_{44}(c) - 12 \frac{c}{a} \left( 1 - \frac{c^2}{a^2} \right) \left\{ \sin^{-1} \frac{a}{c} - \right. \\
 & - \left. \frac{a}{c} \sqrt{1 - \frac{a^2}{c^2}} \right\} \psi_{15}(c) + \frac{c^3}{a^3} \left\{ \sin^{-1} \frac{a}{c} - \frac{a}{c} \sqrt{1 - \frac{a^2}{c^2}} \left( 1 + \frac{4}{3} \frac{a^2}{c^2} - \frac{10}{3} \frac{a^4}{c^4} \right) \right\} \psi_{35}(c) + \\
 & + \frac{1}{135} \left( 1 - \frac{a^2}{c^2} \right)^{\frac{1}{2}} \left( 1 - 17 \frac{a^2}{c^2} + 16 \frac{a^4}{c^4} \right) \psi_{55}(c) \dots\dots\dots(5.11)
 \end{aligned}$$

$$\begin{aligned}
 & - \frac{945\pi}{32} \psi_{00}(a) + \frac{23625\pi}{512} \psi_{02}(a) + \frac{39375\pi}{1024} \psi_{22}(a) - \frac{76545\pi}{4096} \psi_{04}(a) - \frac{14175\pi}{4096} \psi_{24} \\
 & (a) - \frac{99225\pi}{32768} \psi_{44}(a) + \frac{\pi}{2} \psi_{55}(a) = - \frac{945\pi}{32} \psi_{00}(c) + 648 \frac{c^5}{a^5} \left\{ \left( 2 - 3 \frac{a^2}{c^2} + \right. \right.
 \end{aligned}$$

$$\begin{aligned}
 & + \frac{a^4}{c^4} \sin^{-1} \frac{a}{c} - \frac{a}{c} \left(1 - \frac{a^2}{c^2}\right)^{\frac{1}{2}} \left(2 - \frac{5}{3} \frac{a^2}{c^2} - \frac{1}{3} \frac{a^4}{c^4}\right) \} \psi_{11}(c) + \\
 & + \frac{23625\pi}{512} \psi_{02}(c) - \frac{1575\pi}{1024} \left(21 - 46 \frac{a^2}{c^2}\right) \psi_{22}(c) - 168 \frac{c^5}{a^5} \left\{ \left(2 - \right. \right. \\
 & - 3 \frac{a^2}{c^2} + \frac{a^4}{c^4} \left. \right\} \sin^{-1} \frac{a}{c} - \frac{a}{c} \left(1 - \frac{a^2}{c^2}\right)^{\frac{1}{2}} \left(2 - \frac{5}{3} \frac{a^2}{c^2} - \frac{1}{3} \frac{a^4}{c^4}\right) \} \psi_{13}(c) - \\
 & - 56 \frac{c^5}{a^5} \left\{ 3 \left(1 - \frac{a^2}{c^2}\right) \sin^{-1} \frac{a}{c} - \frac{a}{c} \left(1 - \frac{a^2}{c^2}\right)^{\frac{1}{2}} \left(3 - \frac{a^2}{c^2} - 2 \frac{a^6}{c^6}\right) \right\} \psi_{33}(c) - \\
 & - \frac{76545\pi}{4096} \psi_{04}(c) + \frac{567\pi}{4096} \left(21 - 46 \frac{a^2}{c^2}\right) \psi_{24}(c) - \frac{567\pi}{32768} \left(63 - 552 \frac{a^2}{c^2} + 664 \frac{a^4}{c^4}\right) \psi_{44}(c) \\
 & + 24 \frac{c^5}{a^5} \left\{ \left(2 - 3 \frac{a^2}{c^2} + \frac{a^4}{c^4}\right) \sin^{-1} \frac{a}{c} - \frac{a}{c} \left(1 - \frac{a^2}{c^2}\right)^{\frac{1}{2}} \left(2 - \frac{5}{3} \frac{a^2}{c^2} - \frac{1}{3} \frac{a^4}{c^4}\right) \right\} \psi_{15}(c) + \\
 & + 2 \frac{c^5}{a^5} \left\{ 3 \left(1 - \frac{a^2}{c^2}\right) \sin^{-1} \frac{a}{c} - \frac{a}{c} \sqrt{1 - \frac{a^2}{c^2}} \left(3 - \frac{a^2}{c^2} - 2 \frac{a^6}{c^6}\right) \right\} \psi_{35}(c) + \\
 & + \frac{c^5}{a^5} \left\{ \sin^{-1} \frac{a}{c} - \frac{a}{c} \sqrt{1 - \frac{a^2}{c^2}} \left(1 + \frac{2}{3} \frac{a^2}{c^2} + \frac{2}{15} \frac{a^4}{c^4} + \frac{86}{15} \frac{a^6}{c^6} - \frac{128}{15} \frac{a^8}{c^8}\right) \right\} \psi_{55}(c) \\
 & \dots\dots (5.12)
 \end{aligned}$$

where  $a$  and  $c$  are the inner and outer radii of the air gap.

Those equations can be represented by the matrix equation:

$$X(A_{j-1}^{rj-1} \underline{C}_{j-1} + \underline{S}_{j-1}) = Y(A_{j+1}^{rj} \underline{C}_{j+1} + \underline{S}_{j+1}) \dots\dots (6)$$

where the region  $j$  is considered to be the air gap. It is convenient to include the air gap conditions in the chain (4) as part of the propagation [Ref. 1]. For this purpose we must find fictitious  $A_j^{rj-1}$ ,  $A_j^{rj}$ ,  $\underline{S}_j$  such that the relation

$$A_{j-1}^{rj-1} \underline{C}_{j-1} + \underline{S}_{j-1} = A_j^{rj-1} \underline{C}_j + \underline{S}_j$$

$$A_j^{rj} \underline{C}_j + \underline{S}_j = A_{j+1}^{rj} \underline{C}_{j+1} + \underline{S}_{j+1}$$

is equivalent to the (6). For that we need:

$$\underline{S}_j = 0 \quad A_j^{rj-1} = X^{-1} \quad A_j^{rj} = Y^{-1}$$

## A p p e n d i x 2

### Numerical calculations of the $G(\gamma)$ functions

The formulae given for the  $G(\gamma)$  functions are not suitable for numerical calculations, because, for very small values of  $\alpha$ ,  $\gamma^2 \approx 3\alpha$ , and as it can easily be seen, in the equation (2b) the term  $3\alpha - \gamma^2$  appears.

From the bicubic we can get

$$(3\alpha - \gamma^2) = \frac{1}{385} [25 \gamma^6 - 21 (14 + 11\alpha) \gamma^4 + 1190 \alpha \gamma^2]$$

Having made this substitution we get:

$$G_1(\gamma) = -\frac{\alpha}{\gamma}$$

$$G_2(\gamma) = 0.0324675324 \gamma^4 - 0.381818182 \gamma^2 - 0.3 \gamma^2 \alpha + 1.54545454\alpha$$

$$G_3(\gamma) = -0.041322314 \gamma^5 + 0.381818182 \alpha \gamma^3 + 0.431837856 \gamma^3 - 1.4669421\alpha$$

$$G_4(\gamma) = 0.03047744 \gamma^6 - 0.31045134 \gamma^4 - 0.28161156 \alpha \gamma^4 + 1.007544412 \alpha \gamma^2$$

$$G_5(\gamma) = -0.0138533847 \gamma^7 + 0.1411142654 \gamma^5 + 0.128005276 \alpha \gamma^5 - 0.4579748\alpha$$



### A p p e n d i x 3

#### Input description

Card	Format	Name	Description
1	E 12.5	EPS1	Accuracy on $\gamma$ coefficients
	E 12.5	EPS2	Accuracy on unit matrix
	I 12	NINV1	N° of iterations in 12x12 matrix version
	I 12	NINV2	N° of iterations in 6x6 matrix in sion
	I 12	IMP	= 0 macroscopic input = N microscopic input with a libr of N materials

If IMP = N then N card 2 are needed

2	E 12.5	TRANSP(I)	Microscopic transport cross secti material I
	E 12.5	ABS(I)	Microscopic absorption cross sect material I
	A 12	TITLE(I)	Name of material I
3	I 12	NREGS	N° of regions ( $\leq 10$ )
	I 12	NX	= 0 no effect = 1 print intermediate data (nor- mally not necessary)

If IMP > 0 then cards 4 and 5 are needed for each region

4	E 12.5	RAD(IN)	Outer radius region IN
	E 12.5	SOURCE(IN)	Source (neutrons/unit volume) re- gion IN
	E 12.5	AMAT	N° of materials in region IN
	E 12.5	APRIN(IN)	N° of intermediate print-out poin region IN

Each card 4 is followed by AMAT card 5.

Card	Format	Name	Description
5	I 12	NMAT	Material N° (given by the order in which the material appears in card 2)
	E 12.5	CONC	Concentration
If IMP = 0 then card 6 is needed for each region			
6	E 12.5	RAD(I)	Outer radius region I
	E 12.5	FMP(I)	Mean free path region I (negative means air gap)
	E 12.5	SOURCE(I)	Source (neutrons/unit volume) region I
	E 12.5	ALPHA(I)	N° absorptions/collision region I
	E 12.5	blank	
	E 12.5	APRIN(I)	N° of intermediate print-out per region I
7	I 12	IQ	= 0 infinite medium = 1 reflective boundary.

PSMAIN

## Appendix 4

## Code listing

```

EPS1      ACCURACY ON GAMMA COEFFICIENTS
EPS2      ACCURACY ON UNIT MATRIX
NINV1     NO. OF ITERATIONS IN 12*12 MATRIX INVERSION
NINV2     NO. OF ITERATIONS IN 6*6 MATRIX INVERSION
NREGS     NO. OF REGIONS
NX        =0 NO EFFECT,      =1 PRINT INTERMEDIATE DATA
RAD(IN)   OUTER RADIUS REGION IN
FMP(IN)   MEAN FREE PATH REGION IN
SOURCE(IN) A NEGATIVE MEAN FREE PATH MEANS AIR GAP
ALPHA(IN) SOURCE (NEUTRONS/UNIT VOLUME) REGION IN
AMASS(IN) NO. ABSORPTIONS/COLLISION REGION IN
APRIN(IN) ATOMIC WEIGHT OF SCATTERER REGION IN (NOT USED NOW)
IMP       NUMBER OF INTERMEDIATE PRINT OUT POINTS REGION IN
          =0 MACROSCOPIC INPUT
          =N MICROSCOPIC INPUT WITH A LIBRARY OF N MATERIALS
TRANSP(I) MICROSCOPIC LIBRARY SIGMA TRANSP. MATERIAL I
ABS(I)    MICROSCOPIC LIBRARY SIGMA ABSORPTION MAT. I
TITLE(I)  MICROSCOPIC LIBRARY NAME MATERIAL I
CONC      CONCENTRATION
AMAT      NO. OF MATERIALS REGION IN
NMAT      MATERIAL NUMBER (IN LIBRARY)

DIMENSION A(12,12,10), B(12,12,10), U(12,6,10),
1 VAN(12,12), U1(12,6), RESP(12,12), V(12,1),
2 RV(6), V1(12), S(12,10), C(12,10),
3 RAD(10), FMP(10), SOURCE(10), ALPHA(10), AMASS(10), APRIN(10)
DIMENSION Z(12,12), Z1(12,12)
DIMENSION Y(1), RADM(1)
DIMENSION EPS1(1), EPS2(1)
DIMENSION TRANSP(50), ABS(50), TITLE(50)
COMMON A,B,U,VAN,U1,RESP,V,RV,V1,S,C,RAD,FMP,SOURCE,ALPHA,AMASS,A
1 IRIN,IR,IRR,IN,IQ,N5,N6
COMMON Z,Z1,Y,RADM,NREGS
COMMON NX
COMMON EPS1,EPS2,NINV1,NINV2
1 FORMAT (6E12.5)
2 FORMAT (6I12)
4 FORMAT (2E12.5,3I12)
3001 FORMAT (2E12.5,A12)
3002 FORMAT (1H1,36X,20H MICROSCOPIC LIBRARY/////33X,10H TRANSPORT,11X,
11H ABSORPTION,11X,5H NAME///)
3003 FORMAT (E44.5,E21.5,A17)
3004 FORMAT (12CH1 REGION OUTER RADIUS SOURCE NO. MATE
1IALS NAME CONCENTRATION NO. ABS/COLL MEAN FREE PATH
3005 FORMAT (4E12.5)
3006 FORMAT (18,E18.5,2E16.5)
3007 FORMAT (112,E12.5)
3008 FORMAT (A71,E16.5)
3009 FORMAT (E103.5,E15.5)

TAPE DESIGNATION
N5=5

```



```

P5MAIN
      N6=6
C
      READ INPUT TAPE N5,4,EPS1,EPS2,NINV1,NINV2,IMP
      IF(IMP) 999,999,8000
8000  READ INPUT TAPE N5,8001,(TRANSP(1),ABS(1),TITLE(1),I=1,IMP)
      WRITE OUTPUT TAPE N6,8002
      WRITE OUTPUT TAPE N6,8003,(TRANSP(1),ABS(1),TITLE(1),I=1,IMP)
      999  READ INPUT TAPE N5,2,NREGS,NX
      IF(IMP) 8060,8060,8051
8051  WRITE OUTPUT TAPE N6,8004
      DO 8101 IN=1,NREGS
      READ INPUT TAPE N5,8005,RAD(IN),SOURCE(IN),AMAT,APRIN(IN)
      WRITE OUTPUT TAPE N6,8006,IN,RAD(IN),SOURCE(IN),AMAT
      MAT=AMAT+C.01
      SIGT=0.0
      SIGA=0.0
      DO 8100 I=1,MAT
      READ INPUT TAPE N5,8007,NMAT,CONC
      WRITE OUTPUT TAPE N6,8008,TITLE(NMAT),CONC
      SIGT=SIGT+CONC*TRANSP(NMAT)
8100  SIGA=SIGA+CONC*ABS(NMAT)
      ALPHA(IN)=SIGA/SIGT
      FMP(IN)=1.0/SIGT
      WRITE OUTPUT TAPE N6,8009,ALPHA(IN),FMP(IN)
8101  CONTINUE
      GO TO 8052
8060  READ INPUT TAPE N5,1,(RAD(1),FMP(1),SOURCE(1),ALPHA(1),AMASS(1),
1 APRIN(1),I=1,NREGS)
      WRITE OUTPUT TAPE N6,3,NREGS
      3  FORMAT(11H1,15,9H REGIONS //)
      WRITE OUTPUT TAPE N6,5003,(1,RAD(1),FMP(1),SOURCE(1),ALPHA(1),1
1 NREGS)
5003  FORMAT(///18X,10H REGION = 13/,12X,16H OUTER RADIUS = E12.5/
1 10X,18H MEAN FREE PATH = E12.5/,18X,10H SOURCE = E12.5/
2 28H NO.ABSORPTIONS/COLLISION = E12.5/)
C
C      INFINITE MEDIUM      IQ=0
C      REFLECTIVE BOUNDARY IQ=1
C
8052  READ INPUT TAPE N5,2,IQ
      DO 20 I=1,NREGS
      IF(ALPHA(I)-0.00001)21,20,20
D 21  ALPHA(I)=0.00001
      WRITE OUTPUT TAPE N6,22,I
      22  FORMAT(8HOREGION 15,19H IS PURE SCATTERER //)
      20  CONTINUE
      NR=NREGS-1
      DO 150 INN=1,NR
      IN=INN
C
C      TEST ON AIR GAP
C
      IF(FMP(IN))6500,6501,6501
6501  CONTINUE
D      RADM=RAD(IN)/FMP(IN)
D      Y=ALPHA(IN)

```

```

P5MAIN
  CALL MATRIX
  DO 23 I=1,12
  DO 23 J=1,12
23  A(I,J,IN)=Z(I,J)

  TEST ON AIR GAP
  IF(FMP(IN+1))6502,6500,6500
6502 CALL AIRGAP
  GO TO 150
6500 CONTINUE

  RADM=RAD(IN)/FMP(IN+1)
  Y=ALPHA(IN+1)
  CALL MATRIX
  CALL SIMH(Z,Z1,12,NINV1)
  DO 24 I=1,12
  DO 24 J=1,12
24  B(I,J,IN)=Z1(I,J)
150 CONTINUE

  DIRECT MATRICES ARE STORED IN A(I,J,IN) FOR EACH RADIUS IN
  INVERTED MATRICES ARE STORED IN B(I,J,IN) FOR EACH RADIUS IN

  IF(IQ)101,50,101
101  RADM=RAD(NREGS)/FMP(NREGS)
  Y=ALPHA(NREGS)
  CALL MATRIX
  DO 25 I=1,12
  DO 25 J=1,12
25  A(I,J,NREGS)=Z(I,J)
50  DO 100 I=1,NREGS
  S(1,I)=SOURCE(I)*FMP(I)/ALPHA(I)
100  DO 100 J=2,12
  S(J,I)=0.0

  NOW ALL MATRICES ARE PREPARED
  PREPARE FIRST STARTER

  IN=1
  DO 500 I=1,6
  DO 500 J=1,12
500  U(I,J,IN)=0.0
  J=0
  DO 501 I=1,11,2
  J=J+1
501  U(I,J,IN)=1.0

  PROPAGATION OF STARTER
  NITER=0
1001 DO 1000 IN=2,NREGS

  MATRIX MULTIPLICATION  A*U=U1
  DO 502 I=1,12
  DO 502 J=1,6

```

PSMAIN

```

      U1(I,J)=0.0
      DO 502 IS=1,12
502  U1(I,J)=U1(I,J)+A(I,IS,IN-1)*U(IS,J,IN-1)
      MATRIX MULTIPLICATION  B*U1=U
      DO 503 I=1,12
      DO 503 J=1,6
      U(I,J,IN)=0.0
      DO 503 IS=1,12
503  U(I,J,IN)=U(I,J,IN)+B(I,IS,IN-1)*U1(IS,J)
1000 CONTINUE
      IF(IQ)504,505,504
      MATRIX MULTIPLICATION  A*U=U1
504  DO 506 I=1,12
      DO 506 J=1,6
      U1(I,J)=0.0
      DO 506 IS=1,12
506  U1(I,J)=U1(I,J)+A(I,IS,NREGS)*U(IS,J,NREGS)
      N=1
      I=1
      DO 510 K=1,3
      I=I+K
      DO 510 KK=1,K
      DO 507 J=1,6
507  VAN(N,J)=U1(I,J)
      N=N+1
510  I=I+1
      GO TO 400
505  N=1
      I=1
      DO 508 K=1,6
      DO 509 J=1,6
509  VAN(N,J)=U(I,J,NREGS)
      N=N+1
508  I=I+2
      VAN(I,J) IS A 6*6 MATRIX OBTAINED SELECTING FROM THE RESPONSE
      6 ROWS OF INTEREST FOR THE OUTER BOUNDARY CONDITIONS
      CHECK IF VAN(I,J) IS UNIT MATRIX
400  DO 401 I=1,6
      DO 401 J=1,6
      IF(I-J)402,403,402
403  RESP(I,J)=VAN(I,J)-1.0
      GO TO 401
402  RESP(I,J)=VAN(I,J)
401  RESP(I,J)=ABS(RESPI(I,J))
      RMAX=RESP(1,1)
      DO 404 I=1,6
      DO 404 J=1,6
      IF(RMAX-RESP(I,J))405,404,404
405  RMAX=RESP(I,J)

```

```

P5MAIN
  404 CONTINUE
C
  IF(NX)5002,5002,5001
5001 CONTINUE
  DO 5000 KK=1,NREGS
    WRITE OUTPUT TAPE N6,2003,KK
2003 FORMAT(1H0//,17H U MATRIX REGION 15//)
    WRITE OUTPUT TAPE N6,2000,((U(I,J,KK),J=1,6),I=1,12)
2000 FORMAT(6E20.5)
5000 WRITE OUTPUT TAPE N6,2001
2001 FORMAT(1H0//)
    WRITE OUTPUT TAPE N6,2000,((VAN(I,J),J=1,6),I=1,6)
    WRITE OUTPUT TAPE N6,2004,RMAX
2004 FORMAT(7H0 RMAX= E12.5//)
    WRITE OUTPUT TAPE N6,2002
2002 FORMAT(1H1//)
5002 CONTINUE
C
  NITER=NITER+1
  IF(RMAX-EPS2)409,409,407
  407 IF(NITER-10)408,408,409
D 408 CALL SIMH(VAN,RESP,6,NINV2)
C
C   MULTIPLICATION OF OLD STARTER BY THE INVERTED VAN(I,J) AND
C   REPETITION OF THE CYCLE
C
  DO 450 I=1,12
  DO 450 J=1,6
D   U1(I,J)=0.0
  DO 450 IS=1,6
D 450 U1(I,J)=U1(I,J)+U(I,IS,1)*RESP(IS,J)
  DO 451 I=1,12
  DO 451 J=1,6
D 451 U(I,J,1)=U1(I,J)
  GO TO 1001
C
  409 WRITE OUTPUT TAPE N6,410,NITER,RMAX,EPS2
  410 FORMAT(///49H1 MAXIMUM OFF-DIAGONAL TERM OF UNIT MATRIX AFTER
1 18H ITERATIONS RMAX = E12.5,13H WHILE EPS2 = E12.5///)
C
C   PROPAGATION OF A ZERO STARTER TO GET THE EFFECT OF THE SOURCES
C
  406 DO 550 I=1,12
D 550 V(I,1)=0.0
  DO 1010 IN=2,NREGS
  DO 551 I=1,12
D   V1(I)=0.0
  DO 551 IS=1,12
D 551 V1(I)=V1(I)+A(I,IS,IN-1)*V(IS,IN-1)
  DO 552 I=1,12
D 552 V1(I)=V1(I)+S(I,IN-1)-S(I,IN)
  DO 553 I=1,12
D   V(I,IN)=0.0
  DO 553 IS=1,12
D 553 V(I,IN)=V(I,IN)+B(I,IS,IN-1)*V1(IS)
  1010 CONTINUE

```



P5MAIN

```

      IF(1Q)554,555,554
554  DO 560 I=1,12
      V1(I)=0.0
      DO 560 IS=1,12
560  V1(I)=V1(I)+A(I,IS,NREGS)*V(IS,NREGS)
      RV(1)=-V1(2)
      RV(2)=-V1(5)
      RV(3)=-V1(6)
      RV(4)=-V1(10)
      RV(5)=-V1(11)
      RV(6)=-V1(12)
      GO TO 556
555  RV(1)=-V(1,NREGS)
      RV(2)=-V(3,NREGS)
      RV(3)=-V(5,NREGS)
      RV(4)=-V(7,NREGS)
      RV(5)=-V(9,NREGS)
      RV(6)=-V(11,NREGS)

```

CALCULATION OF THE VECTOR C(I,IN) OF THE COEFFICIENTS IN EACH REGION IN

```

556  DO 600 IN=1,NREGS
      DO 650 I=1,12
      C(I,IN)=0.0
      DO 650 IS=1,6
650  C(I,IN)=C(I,IN)+U(I,IS,IN)*RV(IS)
      DO 651 I=1,12
651  C(I,IN)=C(I,IN)+V(I,IN)
600  CONTINUE

      IF(1Q)7000,7001,7000
7001  WRITE OUTPUT TAPE N6,7002
7002  FORMAT(21H CONTROL ON ACCURACY /34H COEFFICIENTS IN THE LAST
      11ION /47H THE COEFFICIENTS MARKED WITH * SHOULD VANISH //)
7003  WRITE OUTPUT TAPE N6,7003,(C(I,NREGS),I=1,12)
7003  FORMAT(E13.5,2H */E13.5)
7000  CONTINUE

      IF(1Q)652,653,652
653  DO 654 I=1,12,2
654  C(I,NREGS)=0.0
652  CONTINUE

      WRITE OUTPUT TAPE N6,3000,(RV(I),I=1,6)
3000  FORMAT(42H EFFECT OF THE SOURCES ON A ZERO STARTER //6E20.5//)
      WRITE OUTPUT TAPE N6,3001
3001  FORMAT(31H COEFFICIENTS FOR EACH REGION //)
      IF(NREGS-6)3002,3002,3003
3002  DO 3006 I=1,12
3006  WRITE OUTPUT TAPE N6,2000,(C(I,IN),IN=1,NREGS)
      GO TO 3004
3003  WRITE OUTPUT TAPE N6,2000,((C(I,IN),IN=1,6),I=1,12)
      J=NREGS-6
      DO 3005 I=1,12
3005  WRITE OUTPUT TAPE N6,2000,(C(I,IN),IN=1,J)

```

P5MAIN

```

3004  CONTINUE

      CALL PRINT
      GO TO 999
      END(1,0,0,0,0,0,0,1,0,0,0,0,0,0,0,0)

```

## MATRIX

```

SUBROUTINE MATRIX
  DIMENSION A(12,12,10), B(12,12,10), U(12,6,10),
1 VAN(12,12), U1(12,6), RESP(12,12), V(12,10),
2 RV(6), V1(12), S(12,10), C(12,10),
3 RAD(10), FMP(10), SOURCE(10), ALPHA(10), AMASS(10), APRIN(10)
  DIMENSION Z(12,12), Z1(12,12)
  DIMENSION Y(1), RADM(1)
  DIMENSION EPS1(1), EPS2(1)
  COMMON A,B,U,VAN,U1,RESP,V,RV,V1,S,C,RAD,FMP,SOURCE,ALPHA,AMASS,
1 IRIN,IR,IRR,IN,IQ,N5,N6
  COMMON Z,Z1,Y,RADM,NREGS
  COMMON NX
  COMMON EPS1,EPS2,NINV1,NINV2

  DIMENSION G1(3),G2(3),G3(3),G4(3),G5(3),T(4)
  DIMENSION ARG(6),BSI0(6),BSI1(6),BSK0(6),BSK1(6),BSK2(6),BSK3(6)
1 BSK4(6),BSK5(6),BSI2(6),BSI3(6),BSI4(6),BSI5(6)
  DIMENSION AA(1),BB(1),CC(1),D(1),E(1),F(1),G(1),H(1),W(1),X(1),
1 SS(1),UU(1),P2(1),P3(1)
  COMMON ARG,BSI0,BSI1,BSK0,BSK1,BSK2,BSK3,BSK4,BSK5,BSI2,BSI3,BSI
1 BSI5

  CALCULATION OF GAMMA COEFFICIENTS
  NITER=0
  T(1)=1.0
  T(2)=0.0
  T(3)=1.0
  T(4)=11.76+55.44*Y
  BB=T(4)
2 T(1)=T(2)
  T(2)=T(3)
  T(3)=T(4)
  T(4)=11.76*T(3)+9.24*Y*T(3)-15.4*T(2)-47.6*Y*T(2)+46.2*Y*T(1)
  T(1)=T(4)/T(3)
  CC=1.0-T(1)/BB
  D=ABS(CC)
  BB=T(1)
  NITER=NITER+1
  IF(D-EPS1)3,3,55
55 IF(NITER-55)2,2,56
56 WRITE OUTPUT TAPE N6,57
57 FORMAT(47H0 CONVERGENCE NOT ACHIEVED AFTER 50 ITERATIONS //)
  WRITE OUTPUT TAPE N6,58,D,EPS1
58 FORMAT(5H D= E12.5 ,12H WHILE EPS1= E12.5//)
3 D=T(1)-11.76-9.24*Y
  E=15.4+47.6*Y+D*T(1)
  F=D*D-4.0*E
  G=SQRT(F)
  H=G-D
  T(2)=0.5*H
  T(3)=E/T(2)
  DO 50 I=1,3
  UU=T(1)
  T(1)=SQRT(UU)
  E=UU*UU

```

## MATRIX

```

F=E*T(I)
H=U*T(I)
AA=-Y/T(I)
BB=0.0324675324*E-0.381818182*UU-0.3*UU*Y+1.54545454*Y
CC=-0.041322314*F+0.381818182*Y*H+0.431937956*H-1.466942148*Y*
W=0.03047744*E*UU-0.31045134*E-0.28161156*Y*E+1.007544412*Y*UU
X=-0.0138533847*E*H+0.1411142654*F+0.12800527652*Y*F-0.45797486

```

1H

```

G1(I)=AA
G2(I)=BB
G3(I)=CC
G4(I)=W
G5(I)=X

```

50 CONTINUE

## ARGUMENTS OF BESSEL FUNCTIONS

```

ARG(1)=RADN*T(1)
ARG(2)=RADN*T(2)
ARG(3)=RADN*T(3)
ARG(4)=3.99101531*RADM
ARG(5)=1.43937374*RADM
ARG(6)=3.3166247905*RADM

```

```

DO 1000 I=1,6
BSI0(I)=B0F(ARG(I))
BSI1(I)=B1F(ARG(I))
BSK0(I)=BK0F(ARG(I))
BSK1(I)=BK1F(ARG(I))
BSK2(I)=BSK0(I)+2.0*BSK1(I)/ARG(I)
BSK3(I)=BSK1(I)+4.0*BSK2(I)/ARG(I)
BSK4(I)=BSK2(I)+6.0*BSK3(I)/ARG(I)
BSK5(I)=BSK3(I)+8.0*BSK4(I)/ARG(I)
AA=0.01*BSI0(I)
BB=0.0
SS=15.0

```

```

DO 30 J=1,11
CC=2.0*SS*AA/ARG(I)+BB
BB=AA
AA=CC

```

30 SS=SS-1.0

```

P2=AA
P3=BB
DO 31 J=1,2
CC=2.0*SS*AA/ARG(I)+BB
BB=AA
AA=CC

```

31 SS=SS-1.0

```

CC=4.0*AA/ARG(I)+BB
CC=BSI1(I)/CC
BSI2(I)=AA*CC
BSI3(I)=BB*CC
BSI4(I)=P2*CC
BSI5(I)=P3*CC

```

1000 CONTINUE

## MATRIX ASSEMBLY

## MATRIX

```

DO 500 I=1,12
DO 500 J=1,12
500 Z(I,J)=0.0
Z(1,1)=BSI0(1)
Z(1,2)=BSK0(1)
Z(1,3)=BSI0(2)
Z(1,4)=BSK0(2)
Z(1,5)=BSI0(3)
Z(1,6)=BSK0(3)
Z(2,1)=G1(1)*BSI1(1)
Z(2,2)=G1(1)*BSK1(1)
Z(2,3)=G1(2)*BSI1(2)
Z(2,4)=G1(2)*BSK1(2)
Z(2,5)=G1(3)*BSI1(3)
Z(2,6)=G1(3)*BSK1(3)
Z(3,1)=G2(1)*BSI0(1)
Z(3,2)=G2(1)*BSK0(1)
Z(3,3)=G2(2)*BSI0(2)
Z(3,4)=G2(2)*BSK0(2)
Z(3,5)=G2(3)*BSI0(3)
Z(3,6)=G2(3)*BSK0(3)
Z(3,7)=BSI0(4)
Z(3,8)=BSK0(4)
Z(3,9)=BSI0(5)
Z(3,10)=BSK0(5)
Z(4,1)=G2(1)*BSI2(1)
Z(4,2)=G2(1)*BSK2(1)
Z(4,3)=G2(2)*BSI2(2)
Z(4,4)=G2(2)*BSK2(2)
Z(4,5)=G2(3)*BSI2(3)
Z(4,6)=G2(3)*BSK2(3)
Z(4,7)=G2(4)*BSI2(4)
Z(4,8)=G2(4)*BSK2(4)
Z(4,9)=G2(5)*BSI2(5)
Z(4,10)=G2(5)*BSK2(5)
Z(5,1)=G3(1)*BSI1(1)
Z(5,2)=G3(1)*BSK1(1)
Z(5,3)=G3(2)*BSI1(2)
Z(5,4)=G3(2)*BSK1(2)
Z(5,5)=G3(3)*BSI1(3)
Z(5,6)=G3(3)*BSK1(3)
Z(5,7)=1.252814036*BSI1(4)
Z(5,8)=1.252814036*BSK1(4)
Z(5,9)=3.473732958*BSI1(5)
Z(5,10)=3.473732958*BSK1(5)
Z(6,1)=G3(1)*BSI3(1)
Z(6,2)=G3(1)*BSK3(1)
Z(6,3)=G3(2)*BSI3(2)
Z(6,4)=G3(2)*BSK3(2)
Z(6,5)=G3(3)*BSI3(3)
Z(6,6)=G3(3)*BSK3(3)
Z(6,7)=7.51688421*BSI3(4)
Z(6,8)=7.51688421*BSK3(4)
Z(6,9)=20.84239775*BSI3(5)
Z(6,10)=20.84239775*BSK3(5)
Z(7,1)=G4(1)*BSI0(1)

```



## MATRIX

```

0 Z(7,2)=0.375*G4(1)*BSK2(1)
0 Z(7,3)=0.375*G4(2)*BSK2(2)
0 Z(7,4)=0.375*G4(2)*BSK2(2)
0 Z(7,5)=0.375*G4(3)*BSK2(3)
0 Z(7,6)=0.375*G4(3)*BSK2(3)
0 Z(7,7)=0.233553315*BSK2(4)
0 Z(7,8)=0.233553315*BSK2(4)
0 Z(7,9)=-0.99112907*BSK2(5)
0 Z(7,10)=-0.99112907*BSK2(5)
0 Z(7,11)=BSK2(6)
0 Z(7,12)=BSK2(6)
0 Z(8,1)=-7.5*G4(1)*BSK2(1)
0 Z(8,2)=-7.5*G4(1)*BSK2(1)
0 Z(8,3)=-7.5*G4(2)*BSK2(2)
0 Z(8,4)=-7.5*G4(2)*BSK2(2)
0 Z(8,5)=-7.5*G4(3)*BSK2(3)
0 Z(8,6)=-7.5*G4(3)*BSK2(3)
0 Z(8,7)=-2.89263978*BSK2(4)
0 Z(8,8)=-2.89263978*BSK2(4)
0 Z(8,9)=11.8935439*BSK2(5)
0 Z(8,10)=11.8935439*BSK2(5)
0 Z(8,11)=12.0*BSK2(6)
0 Z(8,12)=12.0*BSK2(6)
0 Z(9,2)=105.0*G4(1)*BSK4(1)
0 Z(9,3)=105.0*G4(2)*BSK4(2)
0 Z(9,4)=105.0*G4(2)*BSK4(2)
0 Z(9,5)=105.0*G4(3)*BSK4(3)
0 Z(9,6)=105.0*G4(3)*BSK4(3)
0 Z(9,7)=-39.236957*BSK4(4)
0 Z(9,8)=-39.236957*BSK4(4)
0 Z(9,9)=166.509685*BSK4(5)
0 Z(9,10)=166.509685*BSK4(5)
0 Z(9,11)=24.0*BSK4(6)
0 Z(9,12)=24.0*BSK4(6)
0 Z(10,1)=1.875*G5(1)*BSK1(1)
0 Z(10,2)=-1.875*G5(1)*BSK1(1)
0 Z(10,3)=1.875*G5(2)*BSK1(2)
0 Z(10,4)=-1.875*G5(2)*BSK1(2)
0 Z(10,5)=1.875*G5(3)*BSK1(3)
0 Z(10,6)=-1.875*G5(3)*BSK1(3)
0 Z(10,7)=-1.7794919942*BSK1(4)
0 Z(10,8)=1.7794919942*BSK1(4)
0 Z(10,9)=2.723518954*BSK1(5)
0 Z(10,10)=-2.723518954*BSK1(5)
0 Z(10,11)=-2.7136021*BSK1(6)
0 Z(10,12)=2.7136021*BSK1(6)
0 Z(11,1)=-52.5*G5(1)*BSK3(1)
0 Z(11,2)=52.5*G5(1)*BSK3(1)
0 Z(11,3)=-52.5*G5(2)*BSK3(2)
0 Z(11,4)=52.5*G5(2)*BSK3(2)
0 Z(11,5)=-52.5*G5(3)*BSK3(3)
0 Z(11,6)=52.5*G5(3)*BSK3(3)
0 Z(11,7)=21.35390393*BSK3(4)
0 Z(11,8)=-21.35390393*BSK3(4)
0 Z(11,9)=-32.6822274*BSK3(5)

```

## MATRIX

```

D      Z(11,10)=32.6822274*BSK3(5)
D      Z(11,11)=-97.689675*BSI3(6)
D      Z(11,12)=97.689675*BSK3(6)
D      Z(12,1)=945.0*G5(1)*BSI5(1)
D      Z(12,2)=-945.0*G5(1)*BSK5(1)
D      Z(12,3)=945.0*G5(2)*BSI5(2)
D      Z(12,4)=-945.0*G5(2)*BSK5(2)
D      Z(12,5)=945.0*G5(3)*BSI5(3)
D      Z(12,6)=-945.0*G5(3)*BSK5(3)
D      Z(12,7)=640.61711781*BSI5(4)
D      Z(12,8)=-640.61711781*BSK5(4)
D      Z(12,9)=-980.4668234*BSI5(5)
D      Z(12,10)=980.4668234*BSK5(5)
D      Z(12,11)=-325.63225*BSI5(6)
D      Z(12,12)=325.63225*BSK5(6)
C
      IF(NX)7000,7001,7000
7000  CONTINUE
      WRITE OUTPUT TAPE N6,3007,T(1),T(2),T(3)
3007  FORMAT(11H1 GAMMA(1)=E12.5/11H GAMMA(2)=E12.5/11H GAMMA(3)=
1  E12.5//)
      WRITE OUTPUT TAPE N6,3001
3001  FORMAT(21HC G COEFFICIENTS // 5X,2HC1,10X,2HC2,10X,2HC3,10X
1  2HC4,10X,2HC5//)
      WRITE OUTPUT TAPE N6,3000,(G1(I),G2(I),G3(I),G4(I),G5(I),I=1,3)
3000  FORMAT(5E12.5)
      WRITE OUTPUT TAPE N6,3002,RADM,Y
3002  FORMAT(8HC RADM= E12.5,5X,8H ALPHA= E12.5)
      WRITE OUTPUT TAPE N6,3003,(ARG(I),I=1,6)
3003  FORMAT(32HC ARGUMENTS OF BESSEL FUNCTIONS // 6E12.5 ///)
      WRITE OUTPUT TAPE N6,3004
3004  FORMAT(19HC BESSEL FUNCTIONS // 5X,2HI0,10X,2HI1,10X,2HI2,10X,
1  2HI3,10X,2HI4,10X,2HI5/)
      WRITE OUTPUT TAPE N6,3005,(BSI0(I),BSI1(I),BSI2(I),BSI3(I),BSI4(
1  BSI5(I),I=1,6)
3005  FORMAT(6E12.5)
      WRITE OUTPUT TAPE N6,3006
3006  FORMAT(2HC 3X,2HK0,10X,2HK1,10X,2HK2,10X,2HK3,
1  10X,2HK4,10X,2HK5/)
      WRITE OUTPUT TAPE N6,3005,(BSK0(I),BSK1(I),BSK2(I),BSK3(I),BSK4(
1  BSK5(I),I=1,6)
C
      M=12
      N=12
2000  FORMAT (6E20.5)
2001  FORMAT(1H1//)
2002  FORMAT(1HC//)
      WRITE OUTPUT TAPE N6,2001
      IF(M-6)10,10,11
10    N1=M
      WRITE OUTPUT TAPE N6,2000,((Z(I,J),I=1,N1),J=1,N)
      RETURN
11    N1=6
      WRITE OUTPUT TAPE N6,2000,((Z(I,J),J=1,N1),I=1,N)
      WRITE OUTPUT TAPE N6,2002
      WRITE OUTPUT TAPE N6,2000,((Z(I,J),J=7,M),I=1,N)

MATRIX
C
7001  CONTINUE
D      RETURN
      END(1,0,0,0,0,0,1,0,0,0,0,0,0,0)

```



## AIRGAP

SUBROUTINE AIRGAP

```

DIMENSION A(12,12,10), B(12,12,10), U(12,6,10),
1 VAN(12,12), U1(12,6), RESP(12,12), V(12,1),
2 RV(6), V1(12), S(12,10), C(12,10),
3 RAD(10), FMP(10), SOURCE(10), ALPHA(10), AMASS(10), APRIN(10)

```

```

DIMENSION Z(12,12), Z1(12,12)

```

```

DIMENSION Y(1), RADM(1)

```

```

DIMENSION EPS1(1), EPS2(1)

```

```

COMMON A,B,U,VAN,U1,RESP,V,RV,V1,S,C,RAD,FMP,SOURCE,ALPHA,AMASS,

```

```

1RIN,IR,IRR,IN,IQ,N5,N6

```

```

COMMON Z,Z1,Y,RADM,NREGS

```

```

COMMON NX

```

```

COMMON EPS1, EPS2, NINV1, NINV2

```

```

DIMENSION AA(1), BB(1), PI(1), CC(1), D(1), E(1), G(1), H(1), UU(1), VV(1)

```

```

1 F(1), H3(1), Z2(1)

```

```

PI=3.1415926535

```

```

AA=RAD(IN)**3

```

```

BB=RAD(IN)**5

```

```

DO 10 I=1,12

```

```

DO 10 J=1,12

```

```

10 Z(1,J)=0.0

```

```

Z(1,2)=RAD(IN)

```

```

Z(2,5)=RAD(IN)

```

```

Z(3,10)=RAD(IN)

```

```

Z(4,2)=-12.0*AA

```

```

Z(4,5)=2.0*AA

```

```

Z(4,6)=AA

```

```

Z(5,2)=-12.0*AA

```

```

Z(5,5)=-28.0*AA

```

```

Z(5,10)=12.0*AA

```

```

Z(5,11)=AA

```

```

Z(6,2)=1296.0*BB

```

```

Z(6,5)=-336.0*BB

```

```

Z(6,6)=-168.0*BB

```

```

Z(6,10)=48.0*BB

```

```

Z(6,11)=6.0*BB

```

```

Z(6,12)=8B

```

```

Z(7,1)=-PI/4.0

```

```

Z(7,2)=PI/2.0

```

```

Z(7,3)=0.15625*PI

```

```

Z(7,4)=-0.078125*PI

```

```

Z(7,7)=0.03515625*PI

```

```

Z(7,8)=-0.00390625*PI

```

```

Z(7,9)=0.00048828125*PI

```

```

Z(8,1)=-0.009375*PI

```

```

Z(8,3)=-0.0203125*PI

```

```

Z(8,4)=0.01953125*PI

```

```

Z(8,5)=0.5*PI

```

```

Z(8,7)=0.507568359*PI

```

```

Z(8,8)=-0.041748046*PI

```

```

Z(8,9)=-0.000274658*PI

```

```

Z(9,1)=-0.05859375*PI

```

```

Z(9,3)=-0.034790039*PI

```

```

Z(9,4)=0.00305175781*PI

```

## AIRGAP

```

D      Z(9,7)=-1.289520263*PI
D      Z(9,8)=0.015106291*PI
D      Z(9,9)=0.006034332275*PI
D      Z(9,10)=0.5*PI
D      Z(10,1)=0.9375*PI
D      Z(10,3)=-1.171875*PI
D      Z(10,4)=-1.7578125*PI
D      Z(10,6)=0.5*PI
D      Z(10,7)=0.197753906*PI
D      Z(10,8)=0.06591796*PI
D      Z(10,9)=-0.04119873*PI
D      Z(11,1)=1.640625*PI
D      Z(11,3)=5.639648437*PI
D      Z(11,4)=-0.769042968*PI
D      Z(11,7)=-10.03611074*PI
D      Z(11,8)=-2.514770597*PI
D      Z(11,9)=0.028839111*PI
D      Z(11,11)=0.5*PI
D      Z(12,1)=-29.53125*PI
D      Z(12,3)=46.14257812*PI
D      Z(12,4)=38.45214843*PI
D      Z(12,7)=-18.68774414*PI
D      Z(12,8)=-3.46669335*PI
D      Z(12,9)=-3.02810668*PI
D      Z(12,12)=0.5*PI
D      DO 11 I=1,12
D      DO 11 J=1,12
D      11 B(I,J,IN)=Z(I,J)
D
D      SOURCE(IN)=0.0
D      AA=RAD(IN)/RAD(IN+1)
D      BB=1.0-AA*AA
D      BB=SQRT(BB)
D      CC=ATANF(AA/BB)
D      C=AA*BB
D      E=CC-C
D      G=AA*AA
D      H=G*AA
D      UU=G*C
D      VV=UU*AA
D      F=RAD(IN+1)**3
D      B3=BB**3
D      ZZ=RAD(IN+1)**5
D      EE=1.-(2./3.)*C
D      DD=(1.-G)**(3./2.)
D      EEE=CC+D
D      A1=1.-1./C
D      DO 22 I=1,12
D      DO 22 J=1,12
D      22 Z(I,J)=0.0
D      Z(1,2)=RAD(IN+1)
D      Z(2,5)=RAD(IN+1)
D      Z(3,10)=RAD(IN+1)
D      Z(4,6)=F
D      Z(4,2)=-12.*F
D      Z(4,5)=2.*F

```



## AIRGAP

```

Z(5,1)=F
Z(5,2)=-12.*F
Z(5,5)=-28.*F
Z(5,10)=12.*F
Z(6,12)=ZZ
Z(6,2)=1296.*ZZ
Z(6,5)=-336.*ZZ
Z(6,10)=48.*ZZ
Z(6,6)=-168.*ZZ
Z(6,11)=6.*ZZ
Z(7,1)=-PI/4.
Z(7,2)=EEE/AA
Z(7,3)=(5.*PI)/32.
Z(7,4)=-((15.*PI)/64.)*EE
Z(7,6)=(7./45.)*DD
Z(7,7)=(9.*PI)/256.
Z(7,8)=-((3.*PI)/256.)*EE
Z(7,9)=-((PI/2048.)*(15.-40.*G+24.*UU))
Z(7,11)=(11./1260.)*DD
Z(7,12)=(11./4200.)*DD*(1.-(8./3.)*G)
Z(8,1)=-((3.*PI)/32.
Z(8,3)=-((105.*PI)/128.
Z(8,4)=-((15.*PI)/256.)*EE
Z(8,5)=EEE/AA
Z(8,6)=-((1./15.)*DD
Z(8,7)=(2079.*PI)/4096.
Z(8,8)=-((513.*PI)/4096.)*EE
Z(8,9)=-((9.*PI)/32768.)*(15.-40.*G+24.*UU)
Z(8,11)=(11./210.)*DD
Z(8,12)=-((11./6300.)*DD*(1.-(8./3.)*G)
Z(9,1)=-((15.*PI)/256.
Z(9,3)=-((1425.*PI)/4096.
Z(9,4)=-((75.*PI)/8192.)*EE
Z(9,7)=-((42255.*PI)/32768.
Z(9,8)=-((1485.*PI)/32768.)*EE
Z(9,9)=-((7.*PI)/262144.)*(15.-40.*G+24.*UU)
Z(9,10)=EEE/AA
Z(9,11)=-((1./28.)*DD
Z(9,12)=(1./2520.)*DD*(1.-(8./3.)*G)
Z(10,1)=(15.*PI)/16.
Z(10,2)=(12.*A1*E)/AA
Z(10,3)=-((75.*PI)/64.
Z(10,4)=-((375.*PI)/128.)*(1./3.-(14./15.)*G)
Z(10,5)=-((2.*A1*E)/AA
Z(10,6)=(CC-D*(1.+(4./3.)*G-(10./3.)*UU))/H
Z(10,7)=(405.*PI)/2048.
Z(10,8)=-((225.*PI)/2048.)*(1./3.-(14./15.)*G)
Z(10,9)=(15.*PI/16384.)*(35.-392.*G+312.*UU)
Z(10,12)=-((11./945.)*BB*(1.-(17.*G)+16.*UU)
Z(11,1)=(105.*PI)/64.
Z(11,2)=(12.*A1*E)/AA
Z(11,3)=(5775.*PI)/1024.
Z(11,4)=-((175.*PI)/2048.)*(5.-14.*G)
Z(11,5)=(28.*A1*E)/AA
Z(11,7)=-((82215.*PI)/8192.
Z(11,8)=-((2289.*PI)/8192.)*(5.-14.*G)

```

## AIRGAP

```

Z(11,9)=-((21.*PI)/65536.)*(35.-390.*G+312.*UU)
Z(11,10)=-((12.*A1*E)/44
Z(11,11)=(CC-D*(1.+(4./3.)*G-(10./3.)*UU))/H
Z(11,12)=(1./135.)*88*(1.-17.*G+16.*UU)
Z(12,1)=-((945.*PI)/32.
Z(12,2)=(648./VV)*(12.-3.*G+UU)*CC-D*(2.-(5./3.)*G-(1./3.)*UU
Z(12,3)=(23625.*PI)/512.
Z(12,4)=-((1575.*PI)/1024.)*(21.-46.*G)
Z(12,5)=-((168./VV)*(12.-3.*G+UU)*CC-D*(2.-(5./3.)*G-(1./3.)*UU
Z(12,6)=-((56./VV)*(3.*(1.-G)*CC-D*(3.-G-2.*H*H))
Z(12,7)=-((76545.*PI)/4096.
Z(12,8)=-((567.*PI)/4096.)*(21.-46.*G)
Z(12,9)=-((567.*PI)/32768.)*(63.-552.*G+664.*UU)
Z(12,10)=(24./VV)*(12.-3.*G+UU)*CC-D*(2.-(5./3.)*G-(1./3.)*UU
Z(12,11)=(2./VV)*(3.*(1.-G)*CC-D*(3.-G-2.*H*H))
Z(12,12)=(CC-D*(1.+(2./3.)*G+(2./15.)*UU+(86./15.)*H*H-(128./1
1*UU*UU))/VV
8200 FORMAT(//('6E20.5))
8201 FORMAT(1H1,52X,17H AIRGAP MATRICES////19H INNER RADIUS)
8301 FORMAT(//19H OUTER RADIUS)
IF(NX) 8202,8202,8203
8203 WRITE OUTPUT TAPE N6,8201
WRITE OUTPUT TAPE N6,8200,((B(I,J,IN),J=1,6),I=1,12)
WRITE OUTPUT TAPE N6,8200,((B(I,J,IN),J=7,12),I=1,12)
WRITE OUTPUT TAPE N6,8301
WRITE OUTPUT TAPE N6,8200,((Z(I,J),J=1,6),I=1,12)
WRITE OUTPUT TAPE N6,8200,((Z(I,J),J=7,12),I=1,12)
8202 CALL SIMH(1,11,12,NINV1)
DO 23 I=1,12
DO 23 J=1,12
23 A(I,J,IN+1)=Z1(I,J)
RETURN
END(1,0,0,0,0,0,1,0,0,0,0,0,0,0,0)

```

PRINT

SUBROUTINE PRINT

FLUXM(IN)            MEAN FLUX REGION IN  
 VOLUME(IN)        VOLUME REGION IN  
 S(1,IN)            SOURCE VECTOR REGION IN  
 ASFLUX            ASYMPTOTIC FLUX  
 ASCURR            ASYMPTOTIC CURRENT

DIMENSION        A(12,12,10),        B(12,12,10),        U(12,6,10),  
 1 VAN(12,12),    U1(12,6),        RESP(12,12),    V(12,10),  
 2 RV(6),        V1(12),        S(12,10),        C(12,10),  
 3 RAD(10), FMP(10), SOURCE(10), ALPHA(10), AMASS(10), APRIN(10)  
 DIMENSION Z(12,12), Z1(12,12)  
 DIMENSION Y(1), RADM(1)  
 DIMENSION EPS1(1), EPS2(1)  
 COMMON A,B,U,VAN,U1,RESP,V,RV,V1,S,C,RAD,FMP,SOURCE,ALPHA,AMASS,  
 1 IRIN,IR,IRR,IN,I0,N5,N6  
 COMMON Z,Z1,Y,RADM,NREGS  
 COMMON NX  
 COMMON EPS1, EPS2, NINV1, NINV2  
 DIMENSION DELTA(1), FLUX(1), RPRIN(1)  
 DIMENSION XN(1), UN(1), VN(1), YN(1), WN(1), ASFLUX(1), ASCURR(1),  
 1 FLUXM(10)  
 DIMENSION ARG(6), BSI1(6), BSI11(6), BSK0(6), BSK1(6), BSK2(6), BSK3(6),  
 1 BSK4(6), BSK5(6), BSI2(6), BSI3(6), BSI4(6), BSI5(6)  
 COMMON ARG,BSI1,BSI11,BSK0,BSK1,BSK2,BSK3,BSK4,BSK5,BSI2,BSI3,BSI  
 1 BSI5  
 DIMENSION VOLUME(10)  
 NX=0  
 2 FORMAT (F12.5,E15.8,3X,E15.8,3X,E15.8)  
 WRITE OUTPUT TAPE N6,1  
 1 FORMAT (10H1    RADIUS 8X,4HFLUX,8X,15HASYMPTOTIC FLUX,20H    ASYMPT  
 1 IC CURRENT //)  
 RPRIN=0.0  
 FLUX=C(1,1)+C(3,1)+C(5,1)+S(1,1)  
 WRITE OUTPUT TAPE N6,2,RPRIN,FLUX  
 DO 10 IN=1,NREGS  
 IF(FMP(IN))20,21,21  
 20 WRITE OUTPUT TAPE N6,22  
 22 FORMAT (11H0    AIR GAP //)  
 RPRIN=RAD(IN)  
 GO TO 10  
 21 N=APRIN(IN)+2.01  
 DELTA=(RAD(IN)-RPRIN)/(APRIN(IN)+1.0)  
 XN=0.0  
 YN=2.0\*FMP(1)/(RAD(1)\*RAD(1))  
 DO 11 IP=1,N  
 IF(RPRIN)11,11,12  
 12 RADM=RPRIN/FMP(IN)  
 Y=ALPHA(IN)  
 CALL MATRIX  
 IF(IP-1)51,50,51  
 51 CONTINUE

## PRINT

```

C      MATRIX MULTIPLICATION AND ADDITION  V1=A*C+S
D      DO 13 I=1,12
D      V1(I)=0.0
D      DO 13 IS=1,12
D 13   V1(I)=V1(I)+Z(I,IS)*C(IS,IN)
D      DO 14 I=1,12
D 14   V1(I)=V1(I)+S(I,IN)
D      FLUX=V1(1)
D      ASFLUX=Z(1,5)*C(5,IN)+Z(1,6)*C(6,IN)+S(1,IN)
D      ASCURR=Z(2,5)*C(5,IN)+Z(2,6)*C(6,IN)
D      WRITE OUTPUT TAPE N6,2,RPRIN,FLUX,ASFLUX,ASCURR
D      IF(IP-N)52,53,52
D 52   CONTINUE
D 11   RPRIN=RPRIN+DELTA
D      RPRIN=RPRIN-DELTA
D 10   CONTINUE
D      IF(IQ)1002,1003,1002
1002  WRITE OUTPUT TAPE N6,1000
1000  FORMAT(21H1 CONTROL ON ACCURACY/29H FLUX AT REFLECTIVE BOUNDA
1 / 45H THE COMPONENTS MARKED WITH * SHOULD VANISH //)
D      WRITE OUTPUT TAPE N6,1001,(V1(I),I=1,12)
1001  FORMAT(E13.5/E13.5,2H */E13.5/E13.5/E13.5,2H */E13.5,2H */
1 E13.5/E13.5/E13.5/E13.5,2H */E13.5,2H */E13.5,2H */)
1003  CONTINUE
D      RETURN
D 50   XN=C(1,IN)*BSI1(1)-C(2,IN)*BSK1(1)
D      XN=RPRIN*XN*RADM/ARG(1)
D      UN=C(3,IN)*BSI1(2)-C(4,IN)*BSK1(2)
D      UN=RPRIN*UN*RADM/ARG(2)
D      VN=C(5,IN)*BSI1(3)-C(6,IN)*BSK1(3)
D      VN=RPRIN*VN*RADM/ARG(3)
D      XN=XN+UN+VN
D      YN=RAD(IN)**2-RPRIN**2
D      YN=2.0*FMP(IN)/YN
D      GO TO 51
D 53   WN=C(1,IN)*BSI1(1)-C(2,IN)*BSK1(1)
D      WN=RPRIN*WN*RADM/ARG(1)
D      UN=C(3,IN)*BSI1(2)-C(4,IN)*BSK1(2)
D      UN=RPRIN*UN*RADM/ARG(2)
D      VN=C(5,IN)*BSI1(3)-C(6,IN)*BSK1(3)
D      VN=RPRIN*VN*RADM/ARG(3)
D      WN=WN+UN+VN-XN
D      FLUXM(IN)=YN*WN+S(1,IN)
D      VOLUME(IN)=3.1415926535*(RAD(IN)*RAD(IN)-RAD(IN-1)*RAD(IN-1))
D      WRITE OUTPUT TAPE N6,60,IN,FLUXM(IN),VOLUME(IN)
60   FORMAT (8H REGION 13,13H MEAN FLUX = E12.5,4X,9HVOLUME = E12.5
D      GO TO 52
D      END(1,0,0,0,0,0,1,0,0,0,0,0,0,0,0)

```

The double precision Bessel  
function subroutines are not  
included in this listing.

3 REGIONS

SAMPLE OUTPUT

```
          REGION = 1  
      OUTER RADIUS = 0.12700E 01  
      MEAN FREE PATH = 0.13828E 01  
          SOURCE = 0.  
NO.ABSORPTIONS/COLLISION = 0.46710E-00
```

```
          REGION = 2  
      OUTER RADIUS = 0.34925E 01  
      MEAN FREE PATH = -0.10000E 01  
          SOURCE = 0.  
NO.ABSORPTIONS/COLLISION = 0.10000E 01
```

```
          REGION = 3  
      OUTER RADIUS = 0.11295E 02  
      MEAN FREE PATH = 0.24612E 01  
          SOURCE = 0.10000E 01  
NO.ABSORPTIONS/COLLISION = 0.10880E-02
```



RADIUS	FLUX	ASYMPTOTIC FLUX	ASYMPTOTIC CURRENT
0.	0.15428335E 03		
0.25400	0.15602420E 03	0.20046557E 03	-0.85674568E 01
0.50800	0.16160909E 03	0.20504069E 03	-0.17330567E 02
0.76200	0.17226632E 03	0.21278194E 03	-0.26489948E 02
1.01600	0.19057240E 03	0.22386651E 03	-0.36256278E 02
1.27000	0.22147449E 03	0.23854876E 03	-0.46855665E 02
REGION 1	MEAN FLUX = 0.18320E 03	VOLUME = 0.50671E 01	

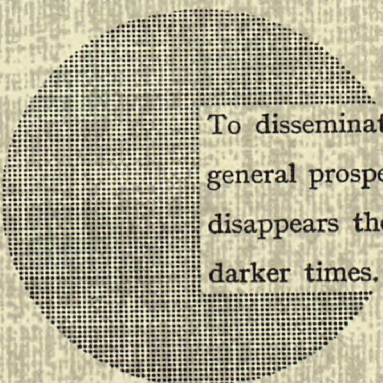
AIR GAP	FLUX	ASYMPTOTIC FLUX	ASYMPTOTIC CURRENT
3.49250	0.26184095E 03	0.26322768E 03	-0.14287844E 02
4.79128	0.28099620E 03	0.28162068E 03	-0.94277818E 01
6.09007	0.29374038E 03	0.29398270E 03	-0.64048030E 01
7.38885	0.30230287E 03	0.30232582E 03	-0.42506678E 01
8.68763	0.30782378E 03	0.30767512E 03	-0.25756345E 01
9.98642	0.31095890E 03	0.31062385E 03	-0.11923563E 01
11.28520	0.31207167E 03	0.31154570E 03	0.54373478E -03
REGION 3	MEAN FLUX = 0.30154E 03	VOLUME = 0.36178E 03	

CONTROL ON ACCURACY  
 FLUX AT REFLECTIVE BOUNDARY  
 THE COMPONENTS MARKED WITH \* SHOULD VANISH

```

0.31207E 03
-0.23668E-10 *
0.40019E-00
-0.32127E 01
0.59039E-06 *
0.31665E-05 *
0.11746E-01
-0.30952E-00
-0.42362E 02
0.79679E-06 *
-0.84697E-05 *
-0.17324E-03 *
```





To disseminate knowledge is to disseminate prosperity — I mean general prosperity and not individual riches — and with prosperity disappears the greater part of the evil which is our heritage from darker times.

Alfred Nobel



## SALES OFFICES

All Euratom reports are on sale at the offices listed below, at the prices given on the back of the cover (when ordering, specify clearly the EUR number and the title of the report, which are shown on the cover).

### PRESSES ACADEMIQUES EUROPEENNES

98, Chaussée de Charleroi, Bruxelles 6

Banque de la Société Générale - Bruxelles  
compte N° 964.558,

Banque Belgo Congolaise - Bruxelles  
compte N° 2444.141,

Compte chèque postal - Bruxelles - N° 167.37,

Belgian American Bank and Trust Company - New York  
compte No. 22.186,

Lloyds Bank (Europe) Ltd. - 10 Moorgate, London E.C.2,

Postcheckkonto - Köln - Nr. 160.861.

### OFFICE CENTRAL DE VENTE DES PUBLICATIONS DES COMMUNAUTES EUROPEENNES

2, place de Metz, Luxembourg (Compte chèque postal N° 191-90)

#### BELGIQUE — BELGIË

MONITEUR BELGE  
40-42, rue de Louvain - Bruxelles  
BELGISCH STAATSBLAD  
Leuvenseweg 40-42 - Brussel

#### GRAND-DUCHE DE LUXEMBOURG

OFFICE CENTRAL DE VENTE  
DES PUBLICATIONS DES  
COMMUNAUTES EUROPEENNES  
9, rue Goethe - Luxembourg

#### DEUTSCHLAND

BUNDESANZEIGER  
Postfach - Köln 1

#### ITALIA

LIBRERIA DELLO STATO  
Piazza G. Verdi, 10 - Roma

#### FRANCE

SERVICE DE VENTE EN FRANCE  
DES PUBLICATIONS DES  
COMMUNAUTES EUROPEENNES  
26, rue Desaix - Paris 15<sup>e</sup>

#### NEDERLAND

STAATSDRUKKERIJ  
Christoffel Plantijnstraat - Den Haag

EURATOM — C.I.D.  
51-53, rue Belliard  
Bruxelles (Belgique)